

MAPPING THE REGION OF INSTABILITY FOR ADIABATIC PACKED BED REACTORS USING A HOMOTOPY CONTINUATION METHOD

DETERMINACIÓN DE LA REGION DE INESTABILIDAD PARA REACTORES ADIABATICOS DE LECHO EMPACADO UTILIZANDO UN METODO DE CONTINUACION POR HOMOTOPIA

JUAN PABLO GUTIÉRREZ-HERNANDEZ

Ingeniero Químico, M. Sc., Universidad Nacional de Colombia, Sede Manizales, J.P.Gutierrez-Hernandez@tue.nl

JAVIER FONTALVO ALZATE

Profesor Asociado, Universidad Nacional de Colombia, Sede Manizales, jfontalvoa@unal.edu.co

MIGUEL ÁNGEL GÓMEZ GARCÍA

Profesor Asociado, Universidad Nacional de Colombia, Sede Manizales, magomez@unal.edu.co

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ABSTRACT: The pioneer schematic ideas of Kimura and Levenspiel (*Ind. Eng. Chem. Proc. Des. Dev.*, 16 (1977) 145 – 148) have been developed to find numerically the region of instability for adiabatic packed bed reactors. Three different cases of special industrial interest and complexity are presented. The highly exothermic gas-phase reactions: ammonia synthesis, methanol production from syn-gas, and SO₂ oxidation. Equations were parameterized and solved according to a continuation homotopy numerical method. The results showed that concentration of inerts and total pressure influences the size of the instability region.

KEYWORDS: Instability region, adiabatic packed bed reactors, homotopy continuation method.

RESUMEN: Las pioneras esquemáticas ideas de Kimura y Levenspiel (*Ind. Eng. Chem. Proc. Des. Dev.*, 16 (1977) 145 – 148) han sido desarrolladas para determinar numéricamente, para reactores de lecho empacado adiabáticos, cuándo un punto de operación es o no estable, y para localizar la envolvente de las condiciones a las cuales se debe evitar su operación. Tres reacciones en fase gas, de especial interés industrial y complejidad, altamente exotérmicas, se utilizan para ilustrar el método de análisis propuesto: la síntesis de amoníaco, la producción de metanol a partir de gas de síntesis, y la oxidación de SO₂. Las ecuaciones son parametrizadas y resueltas de acuerdo a un método de continuación por homotopía. Los resultados demuestran que la concentración de inertes y la presión tienen marcada influencia sobre el tamaño de la región de inestabilidad.

PALABRAS CLAVE: Región de inestabilidad, Reactores de lecho empacado adiabáticos, Método de continuación por homotopía.

1. INTRODUCTION

In the design of adiabatic packed bed reactors, there are two types of temperature excursions that must be specially considered: (1) the existence of “hot spot” temperatures, characteristic of reactors with no recycle of gas; and (2) the unstable “runaway”

temperature, characterized by large recycle of gas. Kimura and Levenspiel [1] schematically have shown that a conversion vs. temperature chart can become a convenient tool for identifying these unstable points. For their recognition, they considered

that “hot spot” and “runaway” temperatures can be related to two ideal reactors models: plug flow tubular reactor (PFTR) and completely stirred tank reactor (CSTR), respectively. With this background, they have developed a criterion to find stable points and an envelope of operating conditions which must be avoided in operation of adiabatic packed bed reactors. From the CSTR model, they linked the molar and adiabatic energy balances in a conversion vs. temperature plot. For a given feed temperature and operating residence time (τ or $\tau' = \tau/C_{A0}$) the adiabatic energy balance line and the S-shaped material balance intersect each other at either one or more stable or unstable points. At the unstable points it is possible to generalize as follows:

$$\left(\begin{array}{l} \text{slope of material} \\ \text{balance curve} \end{array} \right) > \left(\begin{array}{l} \text{slope of adiabatic} \\ \text{energy balance curve} \end{array} \right) \quad (1)$$

Equation (1) was claimed by Kimura and Levenspiel [1] as a fundamental criterion for mapping the instability operation region. Thus, to map the instability region, it is necessary to establish the points where the slope of the constant τ' line (molar balance) equals to the slope of the corresponding adiabatic line (adiabatic energy balance). These points are joined in a τ' -conversion-temperature chart to obtain the boundary of unstable region. Those kinds of charts are used in the literature to solve optimization problems: e.g., to calculate reactor sizes, and to compare alternative designs in reaction rate-conversion-temperature charts [2, 3]. To create those kinds of charts, besides simple rate laws (e.g., first order reactions), computational procedures with complex numerical methods are required due to nonlinearity rate laws and the resulting models. Continuation or homotopy methods have been demonstrated to be robust for the location of all real roots of numerous types of chemical engineering problems involving single or systems of nonlinear equations [4]. Homotopy methods obtain the solution of nonlinear equations by tracking a path from a starting point that is the solution of a much simpler function. Both, the original function and the simpler function are embedded into a parameterized initial-value problem in ordinary differential equations, solved by a predictor–corrector algorithm.

This paper presents how numerically, by using a homotopy continuation method, it is possible to find

the instability regions of complex reaction rate systems, based on the pioneer schematic ideas of Kimura and Levenspiel [1]. Three different cases of special industrial interest and complexity are presented: ammonia, methanol, SO_3 synthesis.

2. MATHEMATICAL MODEL OF UNSTABLE REGIONS IN PACKED BED REACTORS

For the CSTR, the adiabatic energy balance can be arranged as follows [5]:

$$X_A = \frac{\sum_{i=1}^C \theta_i C_{Pi} (T - T_o)}{(-\Delta H_{Rxn})} \quad (2)$$

$$X_A = \frac{\sum_{i=1}^C \theta_i C_{Pi}}{(-\Delta H_{Rxn})} T + \frac{\sum_{i=1}^C \theta_i C_{Pi}}{(-\Delta H_{Rxn})} T_o$$

Let:

$$J = \frac{\sum_{i=1}^C \theta_i C_{Pi}}{(-\Delta H_{Rxn})}$$

and $A = \frac{\sum_{i=1}^C \theta_i C_{Pi}}{(-\Delta H_{Rxn})} T_o$

Then

$$X_A = JT + A \quad (3)$$

where X_A corresponds to the conversion of the limiting reactant. The CSTR molar balance is defined as:

$$\frac{X_A}{-r_A} = \frac{\tau}{C_{A0}} = \tau' \quad (4)$$

The objective is to find the coordinated points of equation (4), (X_A, T) , for a given τ' , with a value of J that satisfy equation (3). It also means to solve the following equation:

$$E(X_A, T) = J - \frac{dX_A}{dT} = 0 \quad (5)$$

with:

$$\frac{dX_A}{dT} = \frac{d}{dT} \left(\frac{\tau(-r_A)}{C_{A0}} \right) = \tau' \frac{d(-r_A)}{dT} \quad (6)$$

It is difficult to evaluate the term $d(-r_A)/dT$ due to its nonlinear dependence on conversion and temperature. This implies to solve the following equation:

$$\left(\frac{dX_A}{dT} \right)_{\tau} = \frac{\tau' \left(\frac{\partial(-r_A)}{\partial T} \right)_{X_A}}{1 - \tau' \left(\frac{\partial(-r_A)}{\partial X_A} \right)_{T}} \quad (7)$$

This is a general expression that corresponds to the instability criteria presented in equation (1). The complexity of applying and solving equation (7) depends on the rate law expression.

To accomplish a numerical solution of equation (1), a second function can be defined as follows:

$$F(X_A, T, \tau') = X_A - \tau'(-r_A) = 0 \quad (8)$$

Thus, equations (5) and (8) have to be solved using a homotopy method (details presented in the appendix). As a result, these two equations are grouped in a new variable named H as follows:

$$H(X_A, T, \tau') = H(Z, \tau') = \begin{bmatrix} E \\ F \end{bmatrix} = 0 \quad (9)$$

Equation (9) contains X_A and T , which are grouped in a new variable Z . As $X_A = X_A(T)$, or $T = T(X_A)$, the derivative of equation (9), with respect to Z is:

$$Z^{k+1} = Z^k - \frac{(\partial H / \partial \tau')}{(\partial H / \partial Z)} \Delta \tau' \quad (10)$$

Basically, the algorithm consists of finding values of the new variable, Z , for a given τ' . This can be made using a multivariable Newton-Raphson method.

3. REACTION SYSTEMS

Here the rate law expressions, for three cases considered in this work, are presented. The rate for ammonia synthesis is [2]:

$$\begin{aligned} N_2 + 3H_2 &\leftrightarrow 2NH_3 \\ r &= k_1 \left(\frac{p_{N_2} p_{H_2}^{1.5}}{p_{NH_3}} \right) - k_2 \left(\frac{p_{NH_3}}{p_{H_2}^{1.5}} \right), \text{ kmol } N_2 \cdot \text{kg}_{\text{cat}}^{-1} \cdot \text{hr}^{-1} \\ k_1 &= 1.79 \times 10^4 \exp\left(-\frac{20800}{RT}\right); \\ k_2 &= 2.57 \times 10^{16} \exp\left(-\frac{47400}{RT}\right) \end{aligned} \quad (11)$$

The formation of methanol based on CO with H_2 is [6]:

$$\begin{aligned} CO + 2H_2 &\leftrightarrow CH_3OH \\ r &= k_1 \left(p_{CO} p_{H_2}^2 - \frac{p_{CH_3OH}}{K_p} \right), \text{ mol } CO \cdot \text{cm}^{-3} \cdot \text{min}^{-1} \\ k_1 &= 2 \times 10^4 \exp\left(-\frac{5000}{T}\right) \\ K_p &= \exp\left(-26.7 + \frac{10913.9}{T}\right) \end{aligned} \quad (12)$$

The oxidation of sulfur dioxide to sulfur trioxide, given by [2]:

$$\begin{aligned} SO_2 + \frac{1}{2}O_2 &\leftrightarrow SO_3 \\ r &= \frac{k_1 p_{O_2} p_{SO_2} \left(1 - \frac{p_{SO_3}}{p_{SO_2} p_{O_2}^{1/2} K_p} \right)}{22.414 (1 + K_2 p_{SO_2} + K_3 p_{SO_3})^2}, \text{ kmol } SO_2 \cdot \text{kg}_{\text{cat}}^{-1} \cdot \text{hr}^{-1} \\ k_1 &= \exp(12.160 - 5473/T) \\ K_2 &= \exp(-9.953 + 8619/T) \\ K_3 &= \exp(-71.745 + 52596/T) \\ K_p &= \exp(11300/T - 10.68) \end{aligned} \quad (13)$$

4. RESULTS

The proposed algorithm proved to be robust and no convergence problems were detected for any of the analyzed rate laws. Figure 1 presents the results obtained for ammonia, methanol and sulfur trioxide synthesis. Calculation time depended on the complexity of the rate law.

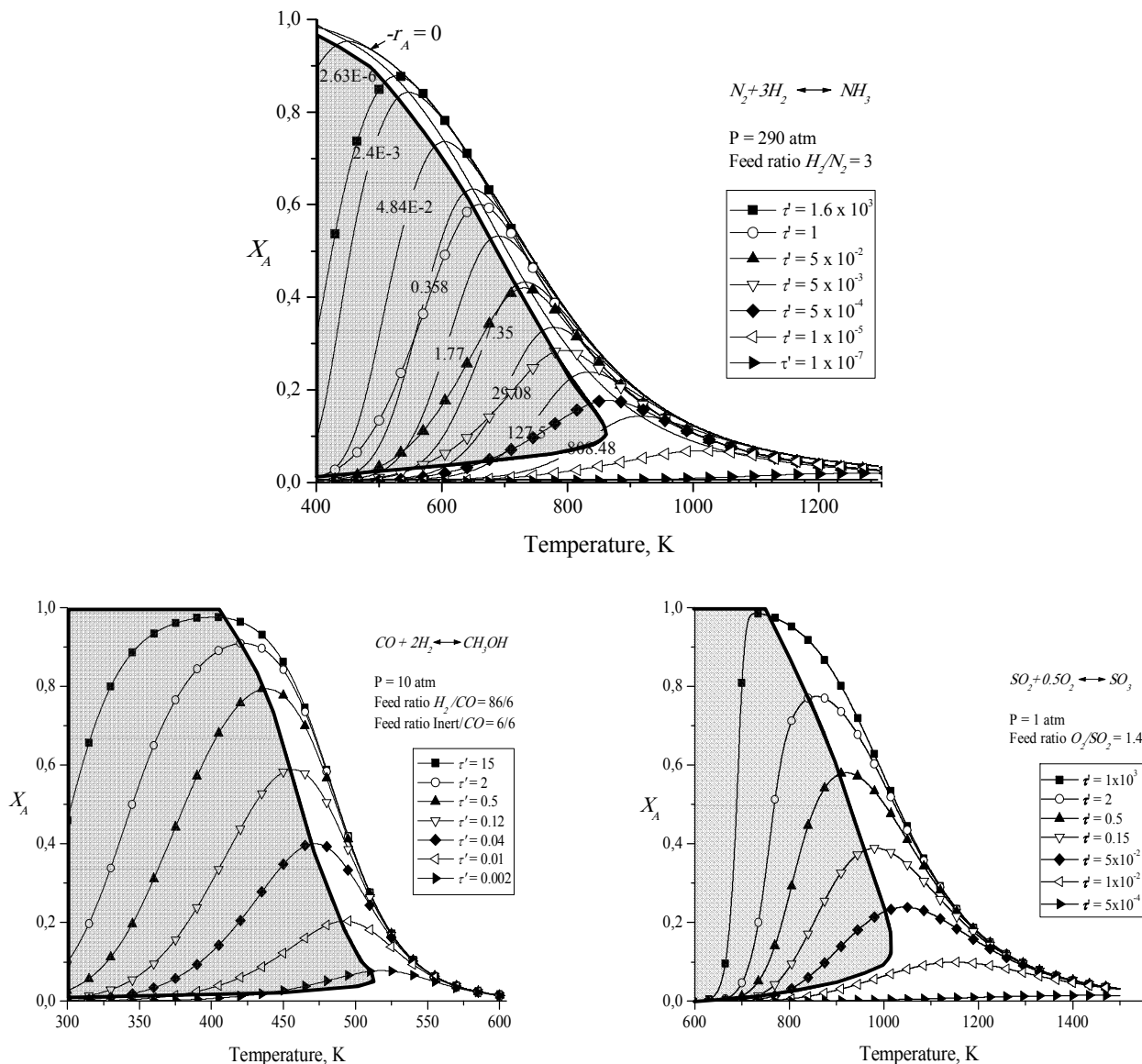


Figure 1. The instability region of the adiabatic packed bed reactors for ammonia, methanol, and SO_3 synthesis. Included curves at τ' constant and, for the ammonia case, also at rate constant

For the three instability regions of Figure 1 computation time were of 115,35; 101,45; and 49,2 seconds for ammonia, methanol, and SO_3 , respectively using a laptop with 1,8 Mhz processor and 776 Mb of RAM. In general, a region of instability is far from the equilibrium ($-r_A = 0$). It is always to the left of the optimal progression of temperature (maximum reaction rate line in the reaction rate – conversion - temperature chart).

The size of the instability region highly depends on the inert concentration as presented for ammonia synthesis (Figure 2). However, for a high concentration of inerts the instability region disappears. Since the instability region has its origins in the difference between generated and removed energy of the system, it can be expected that a decrease in the reactant concentration diminishes the extension of the instability region.

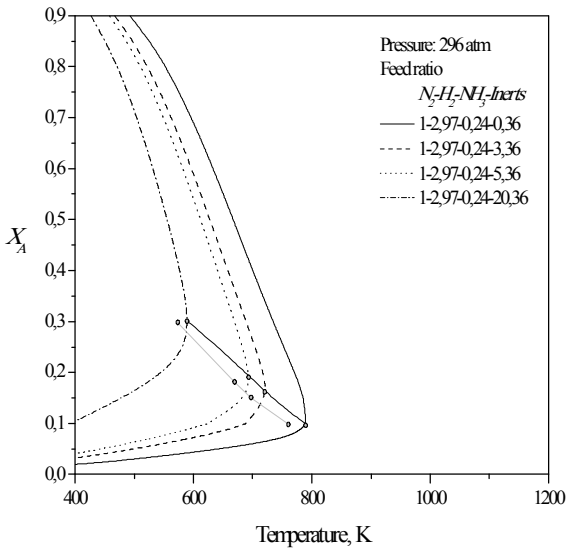


Figure 2. Effect of concentration of inerts on the size of the instability region of an adiabatic packed bed reactor for ammonia synthesis

The trajectories of the instability region have a maximum temperature (Figure 2). The maximum temperature represents the temperature limit for the instability region, and it could be correlated to the adiabatic operation line of a packed bed reactor. The curve of the locus of maxima is also shown in Figure 2. Above this locus curve the derivative dX_A/dT is always negative thus, reactor conversion decreases as temperature increases. Below the locus curve reactor conversions always increases. Consequently, the minimum feed temperature can be found from the adiabatic reactor energy balance (equation (3)). The locus of maxima curve is controlled by operation parameters, such as: type and concentration of inerts (its heat capacity), total pressure, molar feed rate, heat transfer area, temperature of fluid inside the heat-exchange tube; and reactor construction material (its global heat-transfer coefficient) etc... Two of them are presented here: the effect of inert concentration in Figure 2 and the influence of the total pressure on the instability region that is presented in Figure 3, both for ammonia synthesis. In Figure 3, it is possible to observe that upon a decrease on total pressure (at 186 atm, light gray line), the locus curve moves downwards, and slightly to lower temperatures, comparing to those at higher pressures (e.g., 296 atm).

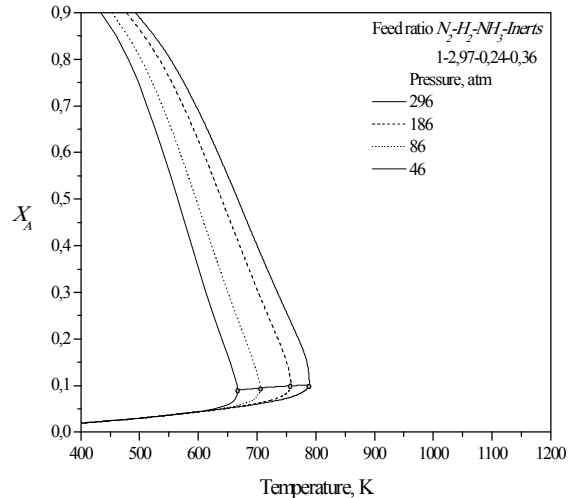


Figure 3. Effect of total pressure on instability region size of an adiabatic packed bed reactor for ammonia synthesis

An increase of total pressure leads to an increase of the extent of the instability region. This can be expected since for an equilibrium reaction, where the number of moles of products is lower than that of reactants, high pressures shifts the equilibrium to the right. Thus, at higher pressures the reaction extent is higher as well as the heat evolution. The locus of maxima of those trajectories is adjusted to straight line, parallel to the temperature axis.

5. CONCLUSIONS

In this paper, a numerical approach, along the original schematic lines of Kimura and Levenspiel [1], for mapping the region of instability of adiabatic packed bed reactors has been shown for three reaction rate laws of high complexity and industrial interest. It was found that the instability region is far from the equilibrium ($-r_A = 0$), and to the left of the optimal progression of temperatures. The shape and size of the instability region of an adiabatic packed bed reactor are controlled by type and concentration of inerts (its heat capacity), total pressure, molar feed rate, heat transfer area, temperature of fluid inside the heat-exchange tube; and reactor construction material. For a high concentration of inerts, the instability region disappears and the increase in total pressure leads to an increase of the extent of the

instability region. The proposed homotopy continuation method appears to be a reasonable simple and accurate numerical technique, for which a simple personal computer amply suffices.

NOMENCLATURE

- C_{AO} = Concentration of reactant in feed, mol/m³.
 C_{pi} = Specific heat of the fluid per mole of entering reactant, cal/(mol.K).
 $-ΔH_{Rxn}$ = Heat of reaction per mole of reactant A, cal/mol.
 k_i = Reaction rate constant.
 K_i = Adsorption equilibrium constant.
 K_p = Partial pressure equilibrium constant.
 p_i = Partial pressure of i .
 R = Ideal gas law constant, 1.9859 cal/(mol.K).
 $-r_A$ = Rate of reaction.
 T = Temperature, K.
 T_o = Entering temperature, K.
 X_A = Fraction of reactant A converted into product.
 $τ$ = Space time based on unit mass or volume of catalyst.
 $θ_i$ = Ratio of the molar flow of species i entering to the molar flow of A entering

APPENDIX

Continuation or homotopy methods have been served as useful tools in modern mathematics and engineering [4]. Stated briefly, a homotopy method consists of the following. Suppose the situation in which very little of *a priori* knowledge concerning zero points of a function $F(x)$ is available. A normal iteration method will often fail to determine the values that make zero the function $F(x)$ because poor starting values are likely to be chosen. As a possible solution, a homotopy or deformation function $H(x,λ)$ is defined. It consists of a linear combination of two functions:

$$F(x) = H(x,1) \quad (A1)$$

$$G(x) = H(x,0) \quad (A2)$$

(A6). It should guaranty that the solution curve Z consists of zero (or near zero) points (u_i) of H .

where $G(x)$ is a function for which a zero is known, and $λ$, the continuation or homotopy parameter, which allows tracking a solution path that connects an arbitrary starting point to the solution of $F(x) = 0$. A standard deformation which is often used is the global homotopy function:

$$H(x,λ) = λF(x) + (1-λ)G(x) = 0 \quad (A3)$$

The choice of $G(x)$ is arbitrary. As the parameter of continuation is in the range [1 0], a series of solutions to $H(x,λ) = 0$ traces a path to the solution of $F(x) = 0$. $H(x,λ)$ is recursively solved at each value of $λ$, using an appropriate numerical method (e.g., a newton type method).

In some instances, even if the function is parametrizable with respect to $λ$, it may be necessary to choose an extremely small increment $Δλ$ for the algorithm to succeed. To remedy this failure or poor performance it can be consider a differential arclength form of equation (A3). Thus, defining a new group of variables:

$$Z(s) = [x(s), λ(s)] \quad (A4)$$

In this way:

$$H(Z(s)) = 0 \quad (A5)$$

Differentiating the equation (5) with respect to arclength s , along the solution path:

$$\frac{\partial H}{\partial Z} \frac{\partial Z}{\partial s} = 0; \quad \left\| \frac{\partial Z}{\partial s} \right\| = 1; \quad (A6)$$

$$Z(0) = (x,1)$$

Defining:

$$\frac{\partial H}{\partial Z} = A, \quad \lambda(A) = \frac{\partial Z}{\partial s} \quad (A7)$$

Where $\lambda(A)$ is the tangent vector induced by A . Now, numerical methods for solving initial value problems may be applied to

Hence, one leads to integrate numerically the equation (A6) very coarsely and then locally use

an iterative method for solving the equation (A5) as a stabilizer.

To describe how the points u_i are generated along the curve Z , it is necessary to suppose that a point u_i has been accepted so that $\|H(u_i)\| < \xi$, where ξ is a numerical value describing the tolerance of the numerical method ($\xi > 0$). This accepted point can be presented as follows:

$$Z_o(s=0) = [x_o, 1] \quad (\text{A8})$$

where x_o is a vector of length defined and with any numerical value.

Such a definition of the equation (A8) is possible since the following statement (A9) is always true:

$$H(X_o, 1) = F(X_o) = 0 \quad (\text{A9})$$

After accepting this point, it proceeds to integrate the following initial value problem:

$$\frac{\partial Z}{\partial s} = \lambda(A), \quad Z(0) = u_i \quad (\text{A10})$$

To obtain a new point u_{i+1} along the curve Z , it is necessary to perform a predictor step due to the difficulty to obtain $\lambda(A)$.

Normally, this predictor step is obtained by a simply numerical integration step of the initial value problem (equation A10). Thus, with an Euler predictor method:

$$v_{i+1} = u_i + \Delta s \cdot \lambda(A), \quad \Delta s > 0 \quad (\text{A11})$$

The point v_{i+1} is just an approximation to the point u_{i+1} . To obtain the correct point that satisfies $\|H(u_{i+1})\| < \xi$, it is necessary to apply a corrector step, such as Newton-type method. This homotopy method applies until $\lambda \leq 0$. This procedure is applied now in two steps for mapping the instability region for tubular packed bed reactors.

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